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SHARP REGULARITY FOR THE INHOMOGENEOUS POROUS MEDIUM EQUATION

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ABSTRACT: We show that locally bounded solutions of the inhomogeneous porous medium equation

$$u_t - \operatorname{div} \left(m u^{m-1} \nabla u \right) = f \in L^{q,r}, \quad m > 1,$$

are locally Hölder continuous, with exponent

$$\gamma = \min\left\{\frac{\alpha_0^-}{m}, \frac{\left[(2q-n)r - 2q\right]}{q\left[(mr - (m-1)\right]}\right\},$$

where α_0 denotes the optimal Hölder exponent for solutions of the homogeneous case. The proof relies on an approximation lemma and geometric iteration in the appropriate intrinsic scaling.

KEYWORDS: Degenerate parabolic equations, porous medium equation, sharp Hölder regularity, intrinsic scaling.

AMS SUBJECT CLASSIFICATION (2010): Primary 35B65. Secondary 35K65, 76S05.

1. Introduction

The quest for obtaining sharp, optimal regularity results is one of the most exciting current trends in the study of nonlinear pdes. Degenerate parabolic equations are known to have Hölder continuous solutions (cf. [9, 16]) under quite general structure assumptions, corresponding to the archetypal p-Laplace equation and porous medium equation (pme). The main difference between these two extensively studied pdes is that the first degenerates at points where the gradient of a solution vanishes and the second at points where this happens for the solution itself. The regularity theory for both equations has evolved in parallel and results for one normally have a counterpart for the other. Recently, in [15], the sharp Hölder exponent

$$\frac{(pq-n)r - pq}{q[(p-1)r - (p-2)]}$$

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for weak solutions of the inhomogeneous *p*-Laplace equation was determined precisely only in terms of p, the space dimension n and the $L^{q,r}$ -integrability of the source. Inspired by the recent breakthroughs in [2, 3, 4], our goal in this paper is to do the same for the porous medium equation (cf. [17]).

Let $U \subset \mathbb{R}^n$ be open and bounded, T > 0 and $U_T = U \times (0,T)$. We consider the prototype inhomogeneous equation

$$u_t - \operatorname{div}\left(mu^{m-1}\nabla u\right) = f, \quad m > 1, \tag{1.1}$$

with a source term $f \in L^{q,r}(U_T) \equiv L^r(0,T;L^q(U))$, where

$$\frac{1}{r} + \frac{n}{2q} < 1, \tag{1.2}$$

which is the standard minimal integrability condition that guarantees the existence of bounded weak solutions and their Hölder regularity.

We will show that bounded weak solutions of (1.1) are locally of class $C^{0,\gamma}$ in space, with

$$\gamma = \frac{\alpha}{m}, \qquad \alpha = \min\left\{\alpha_0^-, \frac{m[(2q-n)r - 2q]}{q[mr - (m-1)]}\right\},\,$$

where $0 < \alpha_0 \leq 1$ denotes the optimal Hölder exponent for solutions of (1.1) with $f \equiv 0$. The regularity class is to be interpreted in the following sense: if

$$\frac{m[(2q-n)r-2q]}{q[mr-(m-1)]} < \alpha_0$$

then solutions are in $C^{0,\gamma}$, with

$$\gamma = \frac{(2q-n)r - 2q}{q[mr - (m-1)]};$$

if, alternatively,

$$\frac{m[(2q-n)r - 2q]}{q[mr - (m-1)]} \ge \alpha_0,$$

then solutions are in $C^{0,\gamma}$, for any $0 < \gamma < \frac{\alpha_0}{m}$. We also obtain the $C^{0,\frac{\gamma}{\theta}}$ regularity in time, where

$$\theta = 2 - \left(1 - \frac{1}{m}\right)\alpha = \alpha \left(1 + \frac{1}{m}\right) + (1 - \alpha) 2$$

is the α -interpolation between $1 + \frac{1}{m}$ and 2. It is worth stressing that, as in the case of the *p*-Laplace equation, the integrability in time (respectively, in space) of the source affects the regularity in space (respectively, in time) of the solution.

We remark that for m = 1 we obtain

$$\gamma = 1 - \left(\frac{2}{r} + \frac{n}{q} - 1\right)$$
 and $\theta = 2$,

recovering the optimal Hölder regularity for the non-homogeneous heat equation, in accordance with estimates obtained by energy considerations.

For n = 1, it is proven in [5] that

$$\alpha_0 = \min\left\{1, \frac{1}{m-1}\right\}$$

but this is not the case in higher dimensions as corroborated by the celebrated counter-example in [6]. The question of the sharp regularity for the homogeneous pme was recently addressed in [13], where it is shown that in the case $m \ge 2$ (see also [14] for 1 < m < 2) a solution achieves the optimal modulus of continuity $C^{0,\frac{1}{m-1}}$ of the Barenblatt fundamental solution after a precise time lag, which is quantified in the paper. This optimal regularity issue is strongly intertwined with the regularity of the free boundary (cf. [8]).

Observe that

$$\frac{m[(2q-n)r-2q]}{q[mr-(m-1)]} = \frac{2m\left(1-\frac{1}{r}-\frac{n}{2q}\right)}{m\left(1-\frac{1}{r}\right)+\frac{1}{r}} > 0$$

and so indeed $\gamma > 0$. Note also that

$$\frac{m[(2q-n)r - 2q]}{q[mr - (m-1)]} > 1$$

if

$$\left(1+\frac{1}{m}\right)\frac{1}{r}+\frac{n}{q}<1,$$

and, as $q, r \to \infty$,

$$\frac{m[(2q-n)r-2q]}{q[mr-(m-1)]} \longrightarrow 2,$$

which means that after a certain integrability threshold it is the optimal regularity exponent of the homogeneous case that prevails, with

$$\alpha = \alpha_0^-$$
 and $\gamma < \frac{\alpha_0}{m} < 1.$

2. Weak solutions and approximation

To fix ideas, we say a nonnegative locally bounded function

$$u \in C_{\text{loc}}(0,T; L^2_{\text{loc}}(U))$$
, with $u^{\frac{m+1}{2}} \in L^2_{\text{loc}}(0,T; W^{1,2}_{\text{loc}}(U))$

is a local weak solution of (1.1) if, for every compact set $K \subset U$ and every subinterval $[t_1, t_2] \subset (0, T]$, we have

$$\int_{K} u\varphi \Big|_{t_1}^{t_2} + \int_{t_1}^{t_2} \int_{K} \left\{ -u\varphi_t + mu^{m-1}\nabla u \cdot \nabla\varphi \right\} = \int_{t_1}^{t_2} \int_{K} f\varphi,$$

for all nonnegative test functions

$$\varphi \in W_{\text{loc}}^{1,2}(0,T;L^2(K)) \cap L^2_{\text{loc}}(0,T;W_0^{1,2}(K)).$$

It is clear that all integrals in the above definition are convergent, interpreting the gradient term as

$$u^{m-1}\nabla u := \frac{2}{m+1}u^{\frac{m-1}{2}}\nabla u^{\frac{m+1}{2}}.$$

An equivalent definition of weak solution involving the Steklov average is instrumental in obtaining the following Caccioppoli estimate.

Proposition 2.1. Let u be a local weak solution of (1.1) and $K \times [t_1, t_2] \subset U \times (0, T]$. There exists a constant C, depending only on n, m and $K \times [t_1, t_2]$, such that

$$\begin{split} \sup_{t_1 < t < t_2} \int_K u^2 \xi^2 + \int_{t_1}^{t_2} \int_K u^{m-1} |\nabla u|^2 \xi^2 \\ &\leq C \int_{t_1}^{t_2} \int_K u^2 \xi \, |\xi_t| + \int_{t_1}^{t_2} \int_K u^{m+1} \left(|\nabla \xi|^2 + \xi^2 \right) + C \|f\|_{L^{q,r}}^2, \\ &\xi \in C_0^\infty(K \times (t_1, t_2)) \text{ such that } \xi \in [0, 1]. \end{split}$$

The proof is standard and follows from testing the equation with $\varphi = u_h \xi^2$, where u_h is the Steklov average of u, and performing the usual combination of integrating in time, passing to the limit in $h \to 0$ and applying Young's inequality (cf. [11, Ch.3, §6]).

We start our fine regularity analysis by fixing the intrinsic geometric setting for our problem. Given $0 < \alpha \leq 1$, let

$$\theta := 2 - \left(1 - \frac{1}{m}\right)\alpha,\tag{2.1}$$

for all

which clearly satisfies the bounds

$$1 + \frac{1}{m} \le \theta < 2.$$

For such θ , define the intrinsic θ -parabolic cylinder as

$$G_{\rho} := \left(-\rho^{\theta}, 0\right) \times B_{\rho}(0), \quad \rho > 0.$$

We next use the available compactness to derive a mechanism linking solutions of the inhomogeneous pme and solutions of the homogeneous equation. This result is to be compared with a similar statement for the *p*-Laplace equation in [15] (see also [12, 7]).

Lemma 2.2. Given $\delta > 0$, there exists $0 < \epsilon \ll 1$ such that if $||f||_{L^{q,r}(G_1)} \leq \epsilon$ and u is a local weak solution of (1.1) in G_1 , with $||u||_{\infty,G_1} \leq 1$, then there exists ϕ such that

$$\phi_t - \operatorname{div}\left(m\phi^{m-1}\nabla\phi\right) = 0 \quad \text{in } G_{1/2} \tag{2.2}$$

and

$$\|u-\phi\|_{\infty,G_{1/2}} \le \delta.$$

Proof: Suppose, for the sake of contradiction, that, for some $\delta_0 > 0$, there exist sequences $(u^j)_j$ and $(f^j)_j$, with

$$u^{j} \in C_{\text{loc}}\left(-1, 0; L^{2}_{\text{loc}}(B_{1})\right), \quad (u^{j})^{\frac{m+1}{2}} \in L^{2}_{\text{loc}}\left(-1, 0; W^{1,2}_{\text{loc}}(B_{1})\right)$$

and $f^j \in L^{q,r}(G_1)$, such that

$$u_t^j - \operatorname{div}\left(m(u^j)^{m-1}\nabla u^j\right) = f^j \quad \text{in } G_1$$
(2.3)

 $\|u^j\|_{\infty,G_1} \le 1,$ (2.4)

$$\|f^{j}\|_{L^{q,r}(G_{1})} \leq 1/j, \qquad (2.5)$$

but still, for any j and any solution ϕ of the homogeneous equation in $G_{1/2}$,

$$\|u^j - \phi\|_{\infty, G_{1/2}} > \delta_0.$$
(2.6)

Consider a cutoff function $\xi \in C_0^{\infty}(G_1)$ such that $\xi \in [0, 1], \xi \equiv 1$ in $G_{1/2}$ and $\xi \equiv 0$ near $\partial_p G_1$. Thus, since u^j is a solution of (1.1), we can apply the Caccioppoli estimate of Proposition 2.1 to get

$$\sup_{-1 < t < 0} \int_{B_1} (u^j)^2 \xi^2 + \int_{-1}^0 \int_{B_1} (u^j)^{m-1} |\nabla u^j|^2 \xi^2$$

$$\leq C \int_{-1}^{0} \int_{B_{1}} (u^{j})^{2} \xi |\xi_{t}| + \int_{-1}^{0} \int_{B_{1}} (u^{j})^{m+1} \left(|\nabla \xi|^{2} + |\xi|^{2} \right) + C ||f^{j}||_{L^{q,r}}^{2} \leq \tilde{c}_{s}$$

using (2.4) and (2.5). Let us now define $v^j := (u^j)^{\frac{m+1}{2}}$. Observing that

$$|\nabla v^j|^2 = \left(\frac{m+1}{2}\right)^2 (u^j)^{m-1} |\nabla u^j|^2,$$

we obtain,

$$\|\nabla v^j\|_{2,G_{1/2}}^2 \le \int_{-1}^0 \int_{B_1} |\nabla v^j|^2 \xi^2 dx dt \le \left(\frac{m+1}{2}\right)^2 \tilde{c}$$

and then, for a subsequence,

$$\nabla v^j \rightharpoonup \psi$$

weakly in $L^2(G_{1/2})$.

Moreover, owing to [10, 9], the equibounded sequence $(u^j)_j$ is also equicontinuous and, by Arzelà–Ascoli theorem,

 $u^j \longrightarrow \phi$

uniformly in $G_{1/2}$, for yet another (relabelled) subsequence. Since we also have the pointwise convergence

$$v^j := (u^j)^{\frac{m+1}{2}} \longrightarrow \phi^{\frac{m+1}{2}} =: v,$$

we can identify $\psi = \nabla v$.

Passing to the limit in (2.3), we find that ϕ solves (2.2) which contradicts (2.6) for $j \gg 1$.

3. Sharp regularity via geometric iteration

We now set up a geometric iteration, exploring the intrinsic scaling of the pme that will be crucial in obtaining the sharp Hölder exponent. The next result is the first step in this iteration. Let α_0 denote the sharp Hölder continuity exponent for solutions of (1.1) in the homogeneous case and

$$\gamma = \frac{\alpha}{m},$$

with

$$0 < \alpha = \min\left\{\alpha_0^-, \frac{m[(2q-n)r - 2q]}{q[mr - (m-1)]}\right\} < \alpha_0 \le \min\left\{1, \frac{1}{m-1}\right\} \le 1. \quad (3.1)$$

Lemma 3.1. There exists $\epsilon > 0$, and $0 < \lambda \ll 1/2$, depending only on m, nand α , such that if $||f||_{L^{q,r}(G_1)} \leq \epsilon$ and u is a local weak solution of (1.1) in G_1 , with $||u||_{\infty,G_1} \leq 1$, then

$$\|u\|_{\infty,G_{\lambda}} \leq \lambda^{\gamma}$$

provided

$$|u(0,0)| \le \frac{1}{4}\lambda^{\gamma}.$$

Proof: Take $0 < \delta < 1$, to be chosen later, and apply Lemma 2.2 to obtain $0 < \epsilon \ll 1$ and a solution ϕ of (2.2) in $G_{1/2}$ such that

$$\|u - \phi\|_{\infty, G_{1/2}} \le \delta$$

Since ϕ solves (2.2), it follows from the available regularity theory (cf. [10]) that ϕ is locally $C_x^{\alpha_0} \cap C_t^{\alpha_0/2}$, for $0 < \alpha_0 < 1$. Thus we obtain

$$\sup_{(x,t)\in G_{\lambda}} |\phi(x,t) - \phi(0,0)| \le C\lambda^{\frac{\alpha_0}{m}},$$

for $\lambda \ll 1$, to be chosen soon, and C > 1 universal. In fact, for $(x, t) \in G_{\lambda}$,

$$\begin{aligned} |\phi(x,t) - \phi(0,0)| &\leq |\phi(x,t) - \phi(0,t)| + |\phi(0,t) - \phi(0,0)| \\ &\leq c_1 |x - 0|^{\alpha_0} + c_2 |t - 0|^{\alpha_0/2} \\ &\leq c_1 \lambda^{\alpha_0} + c_2 \lambda^{\frac{\theta}{2}\alpha_0} \\ &\leq C \lambda^{\frac{\alpha_0}{m}} \end{aligned}$$

since $\theta \ge 1 + \frac{1}{m} > \frac{2}{m}$. We can therefore estimate

$$\sup_{G_{\lambda}} |u| \leq \sup_{G_{1/2}} |u - \phi| + \sup_{G_{\lambda}} |\phi - \phi(0, 0)| + |\phi(0, 0) - u(0, 0)| + |u(0, 0)|$$

$$\leq 2\delta + C\lambda^{\frac{\alpha_0}{m}} + \frac{1}{4}\lambda^{\gamma}.$$
(3.2)

Note that we will choose $\lambda \ll 1/2$ and thus

$$G_{\lambda} := (-\lambda^{\theta}, 0) \times B_{\lambda} \subset (-(1/2)^{\theta}, 0) \times B_{1/2} = G_{1/2}.$$

We finally fix the constants, choosing

$$\lambda = \left(\frac{1}{4C}\right)^{\frac{m}{\alpha_0 - \alpha}}$$
 and $\delta = \frac{1}{4}\lambda^{\gamma}$,

and fixing also $\epsilon > 0$, through Lemma 2.2. The result follows from estimate (3.2) with the indicated choices.

We now iterate the previous result in the appropriate geometric setting.

Theorem 3.2. There exists $\epsilon > 0$, and $0 < \lambda \ll 1/2$, depending only on m, nand α , such that if $||f||_{L^{q,r}(G_1)} \leq \epsilon$ and u is a local weak solution of (1.1) in G_1 , with $||u||_{\infty,G_1} \leq 1$, then

$$\|u\|_{\infty,G_{\lambda^k}} \le (\lambda^k)^{\gamma} \tag{3.3}$$

provided

$$|u(0,0)| \le \frac{1}{4} \left(\lambda^k\right)^{\gamma}.$$

Proof: The proof is by induction on $k \in \mathbb{N}$. If k = 1, (3.3) holds due to Lemma 3.1. Now suppose the conclusion holds for k and let's show it also holds for k + 1. Consider the function $v : G_1 \to \mathbb{R}$ defined by

$$v(x,t) = \frac{u(\lambda^k x, \lambda^{k\theta} t)}{\lambda^{\gamma k}}$$

We have

$$v_t(x,t) = \lambda^{k\theta - \gamma k} u_t(\lambda^k x, \lambda^{k\theta} t),$$

$$\nabla v(x,t) = \lambda^{k-\gamma k} \nabla u(\lambda^k x, \lambda^{k\theta} t)$$

and

$$\operatorname{div}\left(m(v(x,t))^{m-1}\nabla v(x,t)\right)$$

$$= \lambda^{k(2-\alpha)} \operatorname{div} \left(m(u(\lambda^k x, \lambda^{k\theta} t))^{m-1} \nabla u(\lambda^k x, \lambda^{k\theta} t) \right)$$

Recalling (2.1), we conclude, since u is a local weak solution of (1.1) in G_1 , that

$$v_t - \operatorname{div}\left(mv^{m-1}\nabla v\right) = \lambda^{k(2-\alpha)}f(\lambda^k x, \lambda^{k\theta}t) = \tilde{f}(x, t).$$

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We now compute

$$\begin{split} \|\tilde{f}\|_{L^{q,r}(G_{1})}^{r} &= \int_{-1}^{0} \left(\int_{B_{1}} \left| \tilde{f}(x,t) \right|^{q} dx \right)^{r/q} dt \\ &= \int_{-1}^{0} \left(\int_{B_{1}} \lambda^{k(2-\alpha)q} \left| f(\lambda^{k}x,\lambda^{k\theta}t) \right|^{q} dx \right)^{r/q} dt \\ &= \int_{-1}^{0} \left(\int_{B_{\lambda^{k}}} \lambda^{k(2-\alpha)q-kn} \left| f(x,\lambda^{k\theta}t) \right|^{q} dx \right)^{r/q} dt \\ &= \lambda^{[k(2-\alpha)q-kn]\frac{r}{q}} \int_{-1}^{0} \left(\int_{B_{\lambda^{k}}} \left| f(x,\lambda^{k\theta}t) \right|^{q} dx \right)^{r/q} dt \\ &= \lambda^{[k(2-\alpha)q-kn]\frac{r}{q}-k\theta} \int_{-\lambda^{k\theta}}^{0} \left(\int_{B_{\lambda^{k}}} \left| f(x,t) \right|^{q} dx \right)^{r/q} dt. \end{split}$$

Because of the crucial and optimal choice of α in (3.1), we have

$$[k(2-\alpha)q - kn]\frac{r}{q} - k\theta \ge 0$$

and thus

$$\|\tilde{f}\|_{L^{q,r}(G_1)} \le \|f\|_{L^{q,r}((-\lambda^{\theta k}, 0) \times B_{\lambda^k})} \le \|f\|_{L^{q,r}(G_1)} \le \epsilon,$$

which entitles v to Lemma 3.1. Note that $||v||_{\infty,G_1} \leq 1$, due to the induction hypothesis, and

$$|v(0,0)| = \left|\frac{u(0,0)}{(\lambda^k)^{\gamma}}\right| \le \left|\frac{\frac{1}{4}\left(\lambda^{k+1}\right)^{\gamma}}{(\lambda^k)^{\gamma}}\right| \le \frac{1}{4}\lambda^{\gamma}.$$

It then follows that

$$\|v\|_{\infty,G_{\lambda}} \leq \lambda^{\gamma},$$

which is the same as

 $\|u\|_{\infty,G_{\lambda^{k+1}}} \le \lambda^{\gamma(k+1)}.$

The induction is complete.

We next show the smallness regime required in the previous theorem is not restrictive and generalize it to cover the case of any small radius.

Theorem 3.3. If u is a local weak solution of (1.1) in G_1 then, for every $0 < r < \lambda$, we have

$$\|u\|_{\infty,G_r} \le C r^{\gamma}$$

provided

$$|u(0,0)| \le \frac{1}{4}r^{\gamma}.$$

Proof: Take

$$v(x,t) = \rho u\left(\rho^a x, \rho^{(m-1)+2a} t\right)$$

with ρ, a to be fixed, which solves

$$v_t - \operatorname{div}(mv^{m-1}\nabla v) = \rho^{m+2a} f(\rho^a x, \rho^{(m-1)+2a)} t) = \tilde{f}(x, t).$$

We have

$$\|v\|_{\infty,G_1} \le \rho \|u\|_{\infty,G_1}$$

and

$$\|\tilde{f}\|_{L^{q,r}(G_1)}^r = \rho^{(m+2a)r-a(n\frac{r}{q}+2)-(m-1)} \|f\|_{L^{q,r}(G_1)}^r$$

Choosing a > 0 such that

$$(m+2a)r - a\left(\frac{nr}{q}+2\right) - (m-1) > 0,$$

which is always possible, and 0 $< \rho \ll$ 1, we enter the smallness regime required by Theorem 3.2, i.e.,

$$\|v\|_{\infty,G_1} \le 1$$
 and $\|\tilde{f}\|_{L^{q,r}(G_1)} \le \epsilon.$

Now, given $0 < r < \lambda$, there exists $k \in \mathbb{N}$ such that

$$\lambda^{k+1} < r \le \lambda^k.$$

Since $|u(0,0)| \leq \frac{1}{4}r^{\gamma} \leq \frac{1}{4}(\lambda^k)^{\gamma}$, it follows from Theorem 3.2 that

$$\|u\|_{\infty,G_{\lambda^k}} \le (\lambda^k)^{\gamma}.$$

Then, for $C = \lambda^{-\gamma}$,

$$\|u\|_{\infty,G_r} \le \|u\|_{\infty,G_{\lambda^k}} \le (\lambda^k)^{\gamma} < \left(\frac{r}{\lambda}\right)^{\gamma} = C r^{\gamma}$$

We now complete our study, with the main result of the paper.

Theorem 3.4. Let u be a locally bounded weak solution of (1.1) in G_1 , with $f \in L^{q,r}$ satisfying (1.2). Then u is locally of class $C^{0,\gamma}$ in space and $C^{0,\frac{\gamma}{\theta}}$ in time, with

$$\gamma = \frac{\alpha}{m}, \qquad \alpha = \min\left\{\alpha_0^-, \frac{m[(2q-n)r - 2q]}{q[mr - (m-1)]}\right\}.$$

Here $0 < \alpha_0 \leq 1$ denotes the optimal Hölder exponent for solutions of the homogeneous case and θ is given in (2.1).

Proof: We study the Hölder continuity at the origin, proving there is a uniform constant K such that

$$||u - u(0,0)||_{\infty,G_r} \le Kr^{\gamma}.$$
(3.4)

We know, a priori, that u is continuous so we can define

$$\mu := (4|u(0,0)|)^{-\gamma} \ge 0.$$

Take any radius $0 < r < \lambda$. We analyse three alternative cases, exhausting all possibilities.

• If $\mu \leq r < \lambda$ then, by Theorem 3.3,

$$\sup_{G_r} |u(x,t) - u(0,0)| \le C r^{\gamma} + |u(0,0)| \le \left(C + \frac{1}{4}\right) r^{\gamma}.$$
(3.5)

• If $0 < r < \mu$, we consider the function

$$w(x,t) := rac{u(\mu x, \mu^{ heta} t)}{\mu^{\gamma}}.$$

Clearly, $|w(0,0)| = \frac{1}{4}$ and w solves in G_1 the pme

$$w_t - \operatorname{div}\left(mw^{m-1}\nabla w\right) = \mu^{2-\alpha}f(\mu x, \mu^{\theta}t).$$

Moreover, again using Theorem 3.3, it follows that

$$||w||_{\infty,G_1} = \mu^{-\gamma} ||u||_{\infty,G_{\mu}} \le C,$$

since $|u(0,0)| = \frac{1}{4}\mu^{\gamma}$. With this uniform estimate in hand, and using local $C^{0,\alpha}$ regularity estimates [1], we find that there exists a radius ρ_0 , depending only on the data, such that

$$|w(x,t)| \ge \frac{1}{8}, \quad \forall (x,t) \in G_{\rho_0}.$$

This implies that, in G_{ρ_0} , w solves a uniformly parabolic equation of the form

$$w_t - \operatorname{div}\left(a(x,t)\nabla w\right) = f \in L^{q,r},$$

with continuous coefficients satisfying the bounds $0 < c_1 \leq a(x,t) \leq c_2$. In particular, we have (see [15])

$$w \in C^{0,\beta}(G_{\rho_0}), \text{ with } \beta = 1 - \left(\frac{2}{r} + \frac{n}{q} - 1\right) > \gamma,$$

which is the optimal Hölder regularity for solutions of the heat equation with a source in $L^{q,r}$, for exponents satisfying (1.2). As an immediate consequence,

$$\sup_{(x,t)\in G_r} |w(x,t) - w(0,0)| \le C r^{\beta}, \quad \forall \, 0 < r < \frac{\rho_0}{2},$$

which, in terms of u, reads

$$\sup_{(x,t)\in G_r} \left| \frac{u(\mu x, \mu^{\theta} t)}{\mu^{\gamma}} - \frac{u(0,0)}{\mu^{\gamma}} \right| \le C r^{\beta}, \quad \forall \, 0 < r < \frac{\rho_0}{2}.$$

Since $\gamma < \beta$, we conclude

$$\sup_{(x,t)\in G_{\mu r}} |u(x,t) - u(0,0)| \le C \,(\mu r)^{\gamma}, \quad \forall \, 0 < \mu r < \mu \frac{\rho_0}{2},$$

and, relabelling, we obtain

$$\sup_{(x,t)\in G_r} |u(x,t) - u(0,0)| \le C r^{\gamma}, \quad \forall \, 0 < r < \mu \frac{\rho_0}{2}.$$
(3.6)

• Finally, for $\mu \frac{\rho_0}{2} \leq r < \mu$, we have

$$\sup_{(x,t)\in G_{r}} |u(x,t) - u(0,0)| \leq \sup_{(x,t)\in G_{\mu}} |u(x,t) - u(0,0)|$$

$$\leq C \mu^{\gamma} \leq C \left(\frac{2r}{\rho_{0}}\right)^{\gamma} = \tilde{C}r^{\gamma}.$$
(3.7)

Putting $K = \max\left\{C + \frac{1}{4}, \tilde{C}\right\}$ and combining (3.5)–(3.7), we obtain (3.4), for every $0 < r < \lambda$, and the proof is complete.

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