q-Poisson algebras as Poisson-Lie groups

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Abstract

The full set of Poisson-Lie (PL) structures on the 3D book group are presented. They are, by construction, invariant under the group multiplication, and define a multiparameter Poisson-Hopf algebra. We show that the two different *q*-deformed versions of the *sl(2, R)* Poisson algebra are particular cases of Poisson-Lie book groups.

- 1. Motivation **Classifications: state of the art for 3D real Lie groups Poisson-Lie groups** PL groups are quadratic Poisson-Hopf algebras on group manifolds For (semi)simple groups all PL structures are coboundaries. Thus, that are invariant under the group product (the **coproduct map** Δ). classification of classical *r*-matrices (V. Chari, A. Pressley 1994). Quantum Groups are quantizations of PL structures on Lie groups: Heisenberg-Galilei fully classified and obtained therefore, PL groups are the semiclassical limit of quantum groups. (B. Kupershmidt, JPA 1993; A.B., F.J. Herranz, P. Parashar, JPA 1997). PL structures on the group G are in one to one correspondence to Lie Poincaré-Euclidean classification has been also completed bialgebra structures on g = Lie(G) (V.G. Drinfel'd 1986). (P. Stachura, JPA 1998). If the Lie bialgebra (g, δ) is a **coboundary** one, the PL bracket is given All coboundary PL structures obtained for any real group in 3D. by the **Sklyanin bracket** associated to a certain classical *r*-matrix. (A. Rezaei-Aghdam *et al*, JPA 2005).
- Lack of complete classifications for PL structures for non-(semi)simple groups.
- Each PL structure can gives rise to an infinite family of ND integrable **systems** by using the 'coalgebra approach' (A.B., O. Ragnisco 1998).
- Full 3D classification, including many new non-coboundary ones (A.B., A. Blasco, F. Musso, 2011).
- 4D: only oscillator and extended Galilei groups studied (A.B., F.J.) Herranz, JPA 1996; A.B., F.J. Herranz, E. Celeghini, JPA 2000).

2. The book Lie algebra and group

The book Lie algebra (J.F. Cariñena, A. Ibort, G. Marmo, A. Perelomov 1994) is defined by the commutation relations:

$$[e_1, e_3] = e_1$$
 $[e_2, e_3] = e_2$ $[e_1, e_2] = 0$

and is isomorphic to the Bianchi Lie algebra of type V. Lie algebra representation:

$$\rho(e_1) = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad \rho(e_2) = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 0 & 0 \end{pmatrix} \quad \rho(e_3) = \begin{pmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

► The generic Lie group element *M* is given by:

 $M = \exp(z\rho(e_1)) \exp(y\rho(e_2)) \exp(x\rho(e_3)) = \begin{pmatrix} \exp(-x) & 0 & z \\ 0 & \exp(-x) & y \\ 0 & 0 & 1 \end{pmatrix}$

- New variables $X = \exp(-x)$, Y = y, Z = z, then $M = \begin{pmatrix} X & 0 & Z \\ 0 & X & Y \\ 0 & 0 & 1 \end{pmatrix}$
- The multiplication law between two group elements $M \otimes 1$ and $1 \otimes M$ can be read as a **coproduct map** Δ :

 $\Delta(X) = X \otimes X$ $\Delta(Y) = X \otimes Y + Y \otimes 1$ $\Delta(Z) = X \otimes Z + Z \otimes 1$

Equivalently, in 'local' coordinates

$$\Delta(e^{-x}) = e^{-x} \otimes e^{-x} \longrightarrow \Delta(x) = 1 \otimes x + x \otimes 1$$

$$\Delta(y) = e^{-x} \otimes y + y \otimes 1$$

$$\Delta(z) = e^{-x} \otimes z + z \otimes 1$$

3. Classification of PL structures on the book group

Theorem (A.B., A. Blasco, F. Musso, 2011)

The most general Poisson tensor that is quadratic in the group entries of *M* is given by

 $\{X, Y\} = aX^2 - bXY - 2cXZ - aX$ $\{X, Z\} = dX^2 + 2eXY + bXZ - dX$ $\{Y, Z\} = -fX^2 + eY^2 + bYZ - dY + cZ^2 + aZ + f$

where **a**, **b**, **c**, **d**, **e**, **f** are free parameters.

The Casimir function for this Poisson algebra is given by

$$C = \frac{(f(1 + X^2) + d(-1 + X)Y + eY^2 + aZ(1 - X) + Z(bY + cZ))}{X}$$

Coboundary and non-coboundary cases

Most general constant classical r-matrix on the book Lie algebra: $r = r^{12}(e_1 \wedge e_2) + r^{13}(e_1 \wedge e_3) + r^{23}(e_2 \wedge e_3)$ Sklyanin bracket $\{f, g\} = r^{ij} (L_i f L_j g - R_i f R_j g)$ $\{X, Y\} = r^{23}(X-1)X$ $\{X, Z\} = r^{13}(X-1)X$ $\{Y, Z\} = -r^{12}(X^2 - 1) + r^{23}Z - r^{13}Y$

Thus, coboundaries are particular cases with $a = r^{23}, b = 0, c = 0, d = r^{13}, e = 0, f = r^{12}.$ Non-coboundary PL structures: either $b \neq 0$ or $c \neq 0$ or $e \neq 0$.

4. *q*-Poisson algebras are non-coboundary PL structures

The standard *q*-deformation of *sI*(2, *R*)

Change of variables X = e^{-\eta X_0} including a 'deformation' parameter n'

The non-standard q-deformation of sI(2, R)

$$\{X_0, Y\} = \left[a(1 - e^{-\eta X_0}) + bY + 2cZ\right]/\eta$$

$$\{X_0, Z\} = \left[d(1 - e^{-\eta X_0}) - 2eY - bZ\right]/\eta$$

$$\{Y, Z\} = f(1 - e^{-2\eta X_0}) + eY^2 + bYZ - dY + cZ^2 + aZ.$$

If $b = \eta$, $f = 1/\eta$ and a = c = d = e = 0 (non-coboundary!) then

$$\{X_0, Y\} = Y$$

$$\{X_0, Z\} = -Z$$

$$\{Y, Z\} = \frac{1 - e^{-2\eta X_0}}{n} + \eta Y Z$$

which is the standard q-Poisson sl(2)-algebra ($q = e^{\eta}$). \blacktriangleright sl(2) is obtained as the limit $\eta \rightarrow 0$ and the deformed Casimir reads:

 $\mathcal{C}_{\eta} = \frac{1}{n} (e^{\eta X_0} + e^{-\eta X_0}) + \eta Y Z.$

For Through a different change $X = e^{-\eta S_-}$, $Y = S_+$, $Z = S_0$ we get again a multiparametric deformation of a 'classical spin algebra'

$$\{S_{-}, S_{+}\} = \left[a\left(1 - e^{-\eta S_{-}}\right) + bS_{+} + 2cS_{0}\right]/\eta$$

$$\{S_{-}, S_{0}\} = \left[d\left(1 - e^{-\eta S_{-}}\right) - 2eS_{+} - bS_{0}\right]/\eta$$

$$\{S_{+}, S_{0}\} = f\left(1 - e^{-2\eta S_{-}}\right) + eS_{+}^{2} + bS_{+}S_{0} - dS_{+} + cS_{0}^{2} + aS_{0}$$

A non-coboundary case $d = 1, c = -\eta$ and a = b = e = f = 0:

$$egin{aligned} \{m{S}_-,m{S}_+\} &= -2\,m{S}_0\ \{m{S}_-,m{S}_0\} &= rac{1}{\eta}(1-e^{-\eta\,m{S}_-})\ \{m{S}_+,m{S}_0\} &= -m{S}_+ -\eta\,m{S}_0^2 \end{aligned}$$

which is the Poisson version of the **non-standard** $sl_q(2)$ algebra. ► The Casimir function is: $C = \frac{1}{n}(-1 + e^{-\eta S_{-}})S_{+} - e^{\eta S_{-}}S_{0}^{2}$.